# Search for **Hidden Particles**







Particle Physics Seminar — University of Birmingham

March 7, 2018

### Introduction





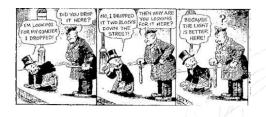


- Higgs found! SM complete and consistent up to Plank scale. But...
  - matter-antimatter asymmetry
  - neutrino masses/mixing
  - dark matter
  - flavour anomalies... New physics?
- NO smoking gun in direct searches up to  $\sim$ 5 TeV...

# What is the energy scale of new physics?

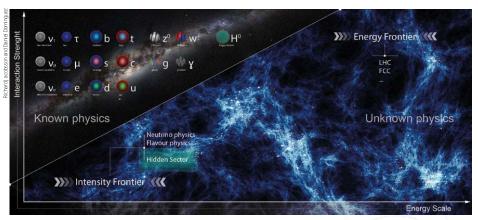


- ightharpoonup Neutrino masses and oscillations: Right Handed see-saw neutrino masses from 1 eV to  $10^{15}~{\rm GeV}$
- $\bullet$  Dark matter: From  $10^{-22}$  eV (super-light scalar) to  $\geq 10^{20}$  GeV (wimpzilla, Q-ball)
- ightharpoonup Baryogenesis: Mass of new particle from  $10~{
  m MeV}$  to  $10^{15}~{
  m GeV}$
- → Higgs mass hierarchy: SUSY, GUT, composite Higgs, large extra dimensions theories require the presence of new particles above the Fermi scale.



# Where is new physics? Experimental approach





- → Unsolved problems ⇒ new particles
- → Why didn't we detect them? Too heavy or too weakly interacting

# Hidden particles



$$\mathcal{L}_{\mathsf{world}} = \mathcal{L}_{\mathsf{SM}} + \mathcal{L}_{\mathsf{portal}} + \mathcal{L}_{\mathsf{HS}}$$

- Hidden Sector (HS) naturally accommodates Dark Matter
  - it may have a rich structure
- Interaction with visible sector (SM) proceeds through mediators
- HS processes very strongly suppressed relative to SM
  - production BRs  $\sim 10^{-10}$
  - very weak interaction with matter
  - very long-lived objects!
- Can search HS through decays to visible particles

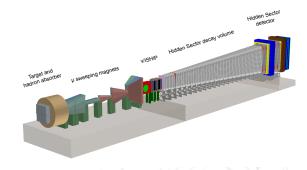
### SHiP: Search for Hidden Particles



SHiP is a new proposed intensity-frontier experiment aiming to search for HNLs and other neutral hidden particles with mass up to  $\mathcal{O}(10)$  GeV and extremely weak couplings.

Aims to be a **zero background** experiment!

Facility also ideally suited for studying  $\nu_{\tau}$  and  $\bar{\nu}_{\tau}$  properties and testing lepton flavour universality by comparing interactions of  $\mu$  and  $\tau$  neutrinos.



### Outline



- → The SHiP experiment
  - Detector system
  - Background strategies
- $\rightarrow$  Physics with  $\nu_{\tau}$
- → The search for Heavy Neutral Leptons
  - Evaluating SHiP sensitivity
- → Probing the Hidden Sector
  - Vector portal
  - Scalar portal
  - Axion-like particles
- → Conclusions

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# Experimental requirements



- → Hidden particles production in **charm** (and beauty) **charm** 
  - Intense proton beam from the SPS (400 GeV)
  - Very dense target of  $10\times\lambda_{\rm int}$ 
    - abundant production of heavy flavours
    - $\bullet$  reduce neutrino production from  $\pi$  and K decays
- → decay of hidden particles:



- large decay volume followed by spectrometer, calorimeter, PID
- shielding from SM particles: hadron absorber + VETO detectors (\*)
- $\rightarrow \tau$  neutrinos:
  - $N_{ au}=4N_{p}\left(\sigma_{c\bar{c}}/\sigma_{pN}\right)f_{D_{s}} imes \mathrm{Br}(D_{s} o au)\simeq\mathbf{6} imes\mathbf{10^{15}}$
  - distinguish  $\nu_{ au}$  /  $\bar{\nu}_{ au}$ : magnetized emulsion target + high-res tracker

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TEASER:  $au 
ightarrow 3 \mu$  sensitivity  $\sim 10^{-10}/\sqrt{N_{
m targets}!}$ 

– distinguish  $\nu_{\tau}$  /  $\bar{\nu}_{\tau}$ : magnetized emulsion target + high-res tracker

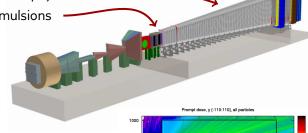
### ...and the muons (\*)?



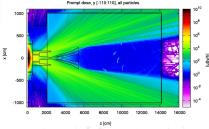
Residual  $\mu$  flux after the hadron absorber is **dangerous**:

background for HS physics

– ageing of  $u_{ au}$  emulsions



- active muon shield based on sweeping magnets
- "conical"-shaped vessel



### History

### 2013:

- submission of the EOI (October, 16 authors)  $\rightarrow arXiv:1310.1762$ 

### 2014:

- SPSC discusses EOI (January)
- 1<sup>st</sup> workshop (June, 100 participants)

### 2015:

- submission of TP (April, 233 authors)
  - $\rightarrow$  arXiv:1504.04956
- submission of PP (April, 85 authors)
  - $\rightarrow$  arXiv:1504.04855
- discussion with SPSC referees

### 2016:

- endorsement by the SPSC (February)

### 2014-today:

12 collaboration meetings, workshops...







### Status and timeline



Accelerator schedule	2015	2016	2017	2018	2019	2020	202	2022	202	23	2024	2025		2026	2027
LHC		F	Run 2			32		Run	3			LS3			Run 4
SPS											NA stop	SPS sto	ор		
					ESPP										
Detector			CD	s	Prototyping,	design		Prod	uction		Insta	llation			
Milestones	TP			CDR		TDR	PRR							CwB	Data taking
Facility						ntegr	ation	-						CwB	
Civil engineering	Pre-construction Target - Detector hall - Beamline - Junction (WP1)														
Infrastructure									nstallation		Installat	tion	Inst.		
Beamline			CD:	S	Prototyping,	design		Production			Install	ation			
Target complex			CD	S	Prototyping,	design	Production Installation								
Target			CD:	s	Prototyping,	design			Produc	ction	In	stallation			

- "Comprehensive Design Study" + Report (2018) to input to the European Strategy consultation in 2019
- New phase of optimisation
  - Re-optimise detector design
  - Assess and improve sensitivity to various hidden sector models

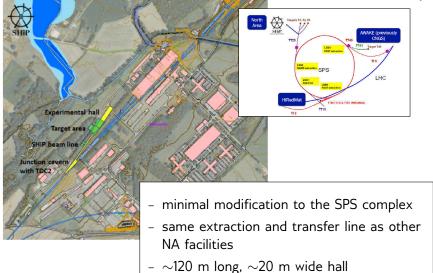
### Outline

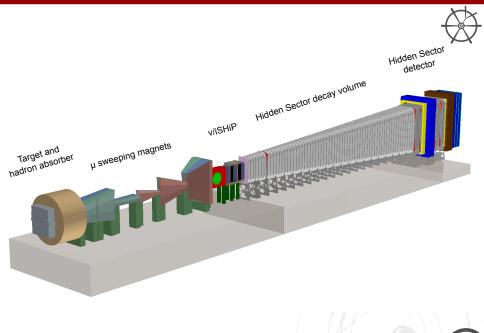


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    - Background strategies
  - Physics with  $u_{ au}$
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# The facility at the SPS

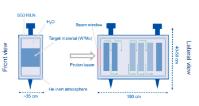






## Target and hadron absorber





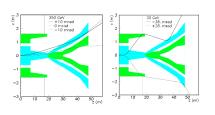


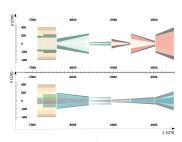
- Layers of titanium/zirconium/molybdenum
  - Heavy target stops hadrons before they decay. Only **muons** emerge
- Peak power 2.5 MW (average 355 kW)
- Challenges:
  - Extremely radioactive environment
  - Many constraints on the integration, access, thermal stress, radioactivity...
- Milestones:
  - Reduced scale prototype by 2019 Q2
  - Full-scale prototype with cooling and power connections by 2020 Q3

### Active muon shield



- Has to reduce  $\mu$  flux by  $\leq 10^{-6}$
- Critical component to optimise to maximize the exp. acceptance
  - global optimisation with machine learning ongoing
- ullet  $\mu$  spectrum measurement planned for 2018 at H4





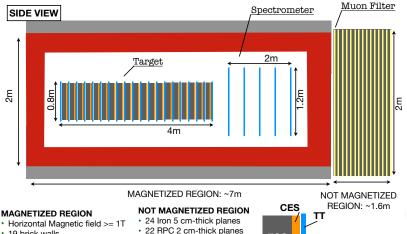
### • Challenges:

- narrow separation between field directions
- aiming to 1.8 T to minimise length
- size of the simulated  $\mu$  sample to optimise with

JINST 12 P05011 2017

### The $\nu$ /iSHiP detector



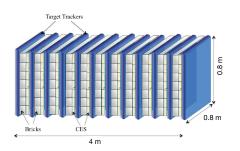


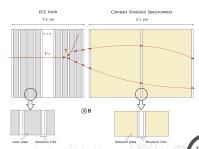
- 19 brick walls
- 20 Target Tracker planes
- 5 Spectrometer planes



# The $\nu$ /iSHiP detector: technologies

- ECC Emulsion Cloud Chambers composed of lead (passive target) and high-res nuclear emulsion films (tracking device)
  - $\longrightarrow \nu$  interaction and  $\tau$  decay vertex with  $\mu \mathrm{m}$  resolution
- CES Compact Emulsion Spectrometer: air gaps and nuclear emulsions
  - $\longrightarrow$  measure the charge of the au daughters, separate  $u_{ au}/ar{
    u}_{ au}$
  - TT Target Trackers: scintillating fibres or MPGD
    - ---- link emulsion tracks to spectrometer tracks
    - ---- provide time stamp and calorimetry information





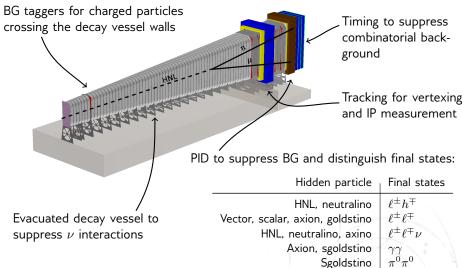
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    - Magnetized tracker for tracking and muon charge: drift tubes
    - Muon filter, also BG tagger for HS detector: iron/RPCs
    - New generation scanning system based on GPU processing

       → scanning 190 cm<sup>2</sup> h<sup>-1</sup>
    - ML techniques for online/offline data processing

### The Hidden Sector detector

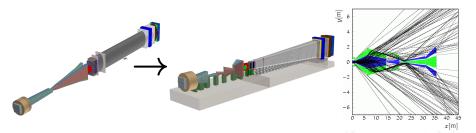




# Hidden Sector detector optimisation



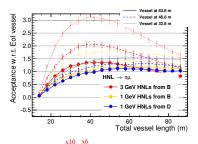
- Optimisation:
  - $\circ$  magnetize hadron stopper  $\longrightarrow$  muon shield shorter and lighter
  - $\circ$  lower  $\mu$  flux entering shield  $\longrightarrow$  decay volume closer to target  $\longrightarrow$  increased signal acceptance and less background
  - decay vessel has now a **pyramidal frustum** shape  $\longrightarrow$  from  $\sim$ 150 m to  $\sim$ 107 m

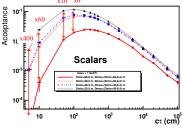


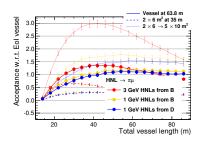
- $\leq 10^{-3}$  mbar to suppress  $\nu$  interactions
- Veto systems surrounding the whole 60 m fiducial volume

# Optimization of the decay volume





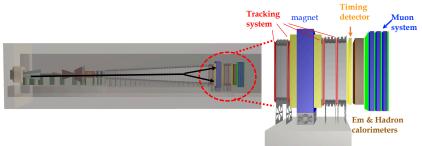




- studied cylindrical, conical solutions in vacuum or He
- might be able to release the vacuum requirement
- → acceptance depends on the hidden particle's lifetime

# Hidden Sector spectrometer





- Full reconstruction and PID of  $e^{\pm}$ ,  $\mu^{\pm}$ ,  $\pi^{\pm}$ ,  $\gamma/\pi^0/\rho^{\pm}$  final states
- NA62-like straw tracker with 5 m long straw tubes
- ullet Timing detector: scintillating bars with SiPM arrays,  $\leq$  100 ps
- Shashlik-type ECAL (HCAL) with 11504 (1512) channels
- Muon detector: 4 xy layers of scintillating bars with WLS fibres (23 km!) and SiPM readout

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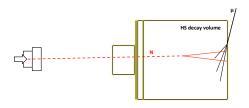
Physics with  $u_{ au}$ 

The search for Heavy Neutral Leptons

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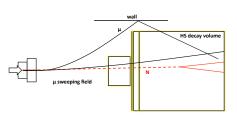


 $\rightarrow$  cosmic  $\mu$  can scatter on the cavern/vessel walls



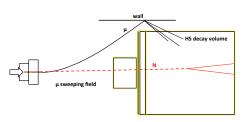


- → cosmic μ can scatter on the cavern/vessel walls
- combinatorial combinations of tracks from different events/vertices



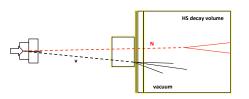


- → cosmic μ can scatter on the cavern/vessel walls
- → combinatorial combinations of tracks from different events/vertices
- μ DIS on the cavern walls can produce charged tracks

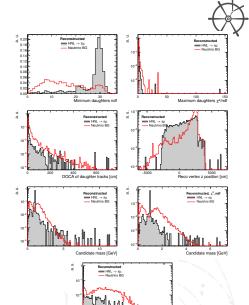




- → cosmic μ can scatter on the cavern/vessel walls
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- ν interactions in the material of the HS detector and upstream closely mimick HP decay topology



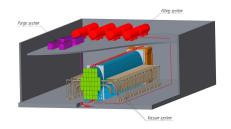
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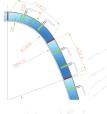


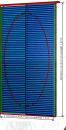
# Shields and background taggers



- → Hadron stopper after the target
- → Magnetic  $\mu$  sweeper creates a  $\geq$  5 m wide fiducial area
- ightharpoonup 
  u detector precedes HS detector and tags upstream particles
- → Upstream VETO complements its acceptance
- → Straw VETO tags decays of  $K_L$  produced in the  $\nu$  detector
- → Liquid scintillator tags interactions crossing the vessel walls
- → Timing detector reduces combinatorial background







### Offline selection



- discard events with activity in the VETO detectors
- select candidates based on the reconstructed direction (must point back to the target)
- require good quality tracks & reconstructed vertex
- event must be fully contained in the fiducial volume, with margins
- we expect < 1 candidate per event  $\rightarrow$  cut on multiplicity

### Selection efficiency

Sample	Multiplicity	Fiducial vol	Track q.	BG cuts/VETO		
$HNL o\pi\mu$	97.5 %	76.1 %	87.0 %	94.2 %		
$\gamma'  o \mu \mu$	99.6 %	85.2 %	94.4 %	94.0 %		
u background	79.1 %	21.0 %	6.5 %	0.0 %		

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Overall  $\lesssim 0.1$  background events / 5 years is attainable!

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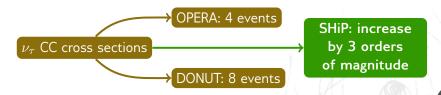
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# $\nu \text{SHiP}$ : a unique opportunity for neutrino physics



High intensity beam dump  $\implies$  high flux of neutrinos (all species).

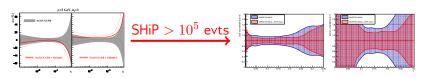
- $\nu_e$  tag: electron shower in the ECC brick
- $\nu_{\mu}$  tag:  $\mu$  in the muon filter
- $u_{ au}$  tag: separation of au production and decay vertices
- first observation of  $\bar{\nu}_{\tau}$
- $\nu_{ au}/\bar{\nu}_{ au}$  cross-section measurements
- ullet First evaluation of the  $F_4$  and  $F_5$  structure functions
- ullet strange quark content from charm production in u scattering



# Tests of perturbative QCD and lepton universality

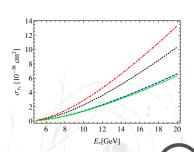


→ PDF improvements with  $\nu$ -nucleon DIS: strange sea quark content currently relies on  $\mathcal{O}(5000)$  charm di- $\mu$  events:



LHC and SHiP will probe different ranges of x.

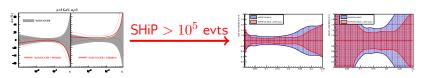
- → Lepton universality tests:
  - hints from LHCb, B factories, ...
  - DIS  $\sigma$  including BSM: Liu, Rashed, Datta PRD92(2015)7, 073016, to compare to  $\sigma_{SM}$
  - results depend on our knowledge of the  $\nu_{ au}$  flux!



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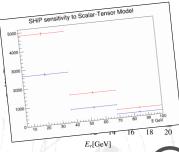


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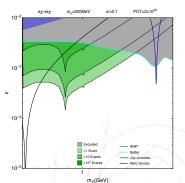


#### iSHiP: Dark matter search



Detect dark matter from dark photon decay through elastic scattering on electrons:  $\chi e^- \to \chi e^-$ . Signature in the emulsion target: a vertex with only  $e^-$  coming out. Simulation  $\Longrightarrow$  background from neutrino scattering can be reduced with kinematical selections to 284 events / 5 y.

Dark photon parameter space for  $\gamma' \to \text{invisible}$  decays excluded by SHiP at 90% C.L., with such expected background and for  $m_\chi = 200$  MeV and  $\chi\gamma'$  coupling  $\alpha' = 0.1$ :



PRD95 (2017) 3, 035006



- The SHiP experiment

  Detector system

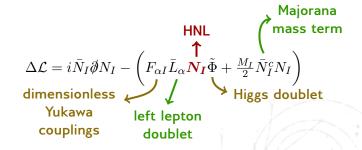
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## Heavy neutral leptons



dark matter neutrino masses/oscillations short-baseline neutrino anomalies matter-antimatter asymmetry

Could be explained with additional, sterile neutrinos



## Heavy Neutral leptons



The Majorana mass term induces 
$$\mathcal{L}_{osc} = c_{\alpha\beta} \left( L_{\alpha}^{\dagger} \tilde{\Phi} \right) \left( \tilde{\Phi} L_{\beta} \right) / \Lambda$$
  $\Longrightarrow$  change flavour of SM neutrino  $\nu_{\alpha} \equiv \tilde{\Phi} L_{\alpha}$ 

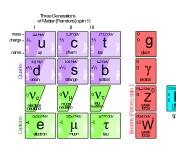
$$m_D = {
m Dirac}$$
 mass term,  $(m_D)_{lpha I} = F_{lpha I} < \Phi >$ 

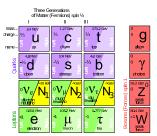
$$\left(\mathcal{M}_{
u}
ight)_{lphaeta} = -\sum_{I} \left(m_{D}
ight)_{lpha I} rac{1}{M_{I}} (m_{D})_{eta I}$$

GeV scale seesaw can generate BAU through HNL oscillations. Because of  $\nu-N$  mixing, HNLs take part in all  $\nu$  processes with strength reduced by  $U_{\alpha I}^2$  and kinematics reflecting  $m_N$ .

# The $\nu$ MSM Asaka, Blanchet, Shaposhnikov, Phys.Lett. B631 (2005) 151-156



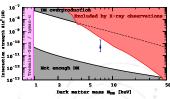






Suitable values of  $m_N$  and  $U_f^2$  allow to simultaneously explain:

- $\nu$  oscillations induced by massive states  $N_2$ ,  $N_3$
- dark matter:  $N_1$  with mass  $\sim \text{keV}$
- BAU: leptogenesis due to Majorana mass term

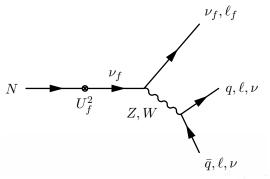


Astrophys. J. 789(2014)13 Phys.Rev.Lett. 113(2014)251301

# HNL phenomenology



HNLs can be produced in decays where a  $\nu$  is replaced by a N (kinetic mixing, low  $\mathcal{BR}$ ). Main neutrino sources in SHiP: c and b mesons.



They can then decay again to SM particles through mixing  $(U^2)$  with a SM neutrino. This (now massive) neutrino can decay to a large amount of final states through emission of a  $Z^0$  or  $W^\pm$  boson.



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# Estimating SHiP's sensitivity to HNLs



→ Number of detected HNL events:

$$\Phi(p.o.t) \times \sigma(pp \to NX) \times \mathcal{P}_{vtx} \times \mathcal{BR}(N \to visible) \times \mathcal{A}$$

with

$$\sigma(pp \to NX) \propto \chi_{cc}, \chi_{bb}, U_f^2$$
  
 $\mathcal{BR}(N \to visible) \propto U_f^2$ 

- → HNL production:
  - $\chi_{cc}$ ,  $\chi_{bb}$  obtained from simulations (Pythia8)
  - $\mathcal{BR}(m_N, U_f^2)$  parametrised according to theory

JHEP 0710 (2007) 015

- $\rightarrow$  Daughters acceptance ( $\mathcal{A}$ ):
  - HNLs kinematics obtained from simulation
  - every decay channel with detectable daughters is simulated

## HNL production in SHiP



Charm mesons are the main source of HNLs in SHiP. Contribution of b mesons for  $m_N>2~{\rm GeV}.$ 

- → Pythia8 used to retrieve the spectrum of c and b mesons in 400 GeV/c proton-on-target collisions
- → HNL production simulated in kinematically-allowed decay chains:
  - $D \to K \ell N$
  - $D_s \rightarrow \ell N$
  - $D_s 
    ightarrow au \, 
    u_ au$  followed by  $au 
    ightarrow \mu \, 
    u \, N$  or  $au 
    ightarrow \pi \, N$
  - $B \to \ell N$
  - $B \to D \ell N$
  - $B_s \to D_s \, \ell \, N$
- $ightharpoonup \mathcal{BR}(pp o NX)$  computed as sum of the BRs of the kinematically-allowed channels

## HNL lifetime and decay channels



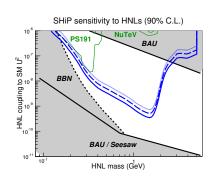
For a given HNL mass, its lifetime was computed on the basis of the widths of its kinematically allowed decay channels:

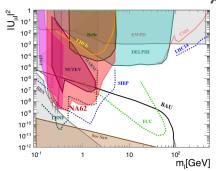
- $N o H^0 
  u$  , with  $H^0 = \pi^0, 
  ho^0, \eta, \eta'$
- $N o H^{\pm} \ell^{\mp}$ , with  $H = \pi, 
  ho$
- $N o 3 \nu$
- $N \to \ell_i^{\pm} \ell_j^{\mp} \nu_j$
- $N \to \nu_i \ell_j^{\pm} \ell_j^{\mp}$
- N o q ar q 
  u and  $N o q ar q \ell$  with q ar q hadronising separately

All decay channels into  $\geq 2$  charged particles considered **detectable**.

# SHiP sensitivity to HNLs







- Best sensitivity up to the charm kinematic limit
- ullet Significant contribution from B decays
- Nice synergy with FCC-ee operated as  $Z^0$  factory  $\circ$  very clean signature  $Z^0 \to \nu \bar{\nu}, \nu \to N \to \ell^+ \ell^-$
- ullet  $10\pm 6$  background events would not affect sensitivity!



- The SHiP experiment
- Detector system
  - Background strategies
- Physics with  $u_{ au}$ 
  - The search for Heavy Neutral Leptons
- Evaluating SHiP sensitivity
- → Probing the Hidden Sector
  - Vector portal
  - Scalar portal
  - Axion-like particles
  - Conclusions

#### Hidden sector



- → new particles are light ⇒ they must be singlets with respect to the gauge group of the SM
- → they may couple to different singlet operators (portals) of the SM
  - dim 2: hypercharge field,  $\epsilon F_{\mu\nu}F'^{\mu\nu}$ , vector portal
    - dim 2: Higgs field,  $\left( lpha_1 \chi + lpha \chi^2 
      ight) H^\dagger H$ , Higgs/scalar portal
    - dim 2 ½: Higgs-lepton,  $YH^T\bar{N}L$ , neutrino portal
    - dim 4:  $AG_{\mu\nu}\epsilon^{\mu\nu\rho\eta}G^{\rho\eta}$ ,  $\partial_{\mu}A\bar{\psi}\gamma^{\mu}\gamma^{5}\psi$ , ..., axion portal
    - SUSY models





The SHiP experiment
Detector system
Background strategies

Physics with  $\nu_{\tau}$ 

The search for Heavy Neutral Leptons

Evaluating SHIP sensitivity

- → Probing the Hidden Sector
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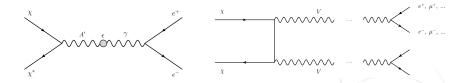
Conclusions

### The vector portal

SM group  $SU(3)\times SU(2)\times U(1)$  may descend from a larger group:

$$SU(3) \times SU(2) \times [U(1)]^n$$





• light ( $m_V \sim \text{MeV-GeV}$ ) vector particles especially interesting for cosmological reasons

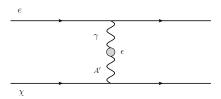
# Dark photons and kinetic mixing



$$\mathcal{L} = \begin{array}{c|c} \hline \textbf{QED-like} \\ \mathcal{L} = \begin{array}{c|c} \mathcal{L}_{\psi,A} & + & \mathcal{L}_{\chi,A'} & -\frac{\epsilon}{2} & F_{\mu\nu}F'^{\mu\nu} & + & \frac{1}{2}m_{A'}^2(A'_{\mu})^2 \\ \hline \uparrow & & \uparrow & & \uparrow \\ \hline \textbf{QED fields} & \textbf{\textit{U}(1)' fields} & & & \text{mass term} \\ \hline \end{array}$$

#### field strength tensors

Eq. of motion:  $\partial_{\mu}F^{\mu\nu}=eJ^{(EM)\nu}\Longrightarrow -\frac{\epsilon}{2}F_{\mu\nu}F'^{\mu\nu}=e\epsilon A'_{\mu}J^{(EM)\mu}$   $\Longrightarrow$  coupling to EM current reduced by  $\epsilon$ .



$$m_{A'} 
ightarrow 0 \Longrightarrow$$
 e.m. charge of  $\chi 
ightarrow e \epsilon.$ 

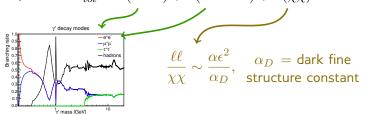
Okun, Sov. Phys. JETP 56 (1982) 502 - Holdom, Phys. Lett. B 166 (1986) 196

## Vector portal phenomenology



→ Decay:

$$\Gamma_{tot} = \Gamma(\ell^+\ell^-) + \Gamma(\text{hadrons}) + \Gamma(\chi\bar{\chi})$$



- → Production at SHiP:
  - meson decays e.g.  $\pi^0 \to \gamma V \; (\sim \epsilon^2)$

Phys.Rev. D80(2009)095024

- p bremsstrahlung on target nuclei pp o ppV

Phys.Lett. B731(2014)320-326

- large  $m_V \Rightarrow$  direct QCD production through underlying  $q\bar{q} \rightarrow V$ ,  $qq \rightarrow V$  (need some more theory work!)

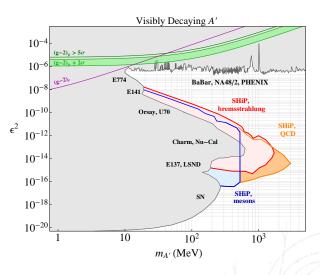
Phys.Rev. D86(2012)035022

→ Light dark matter at SHiP:

if  $\chi \bar{\chi}$  decays dominant  $\Rightarrow \chi$  can scatter on electrons  $\sim \alpha \alpha_D \epsilon^2$ : dense detector to look for light DM.

# SHiP sensitivity: vector portal





Sensitivity studied considering  $\Gamma_{tot} = \Gamma(\ell^+\ell^-) + \Gamma(\mathsf{hadrons})$ .



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## The scalar portal



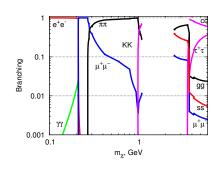
Most general renormalizable  $\mathcal{L}$ :

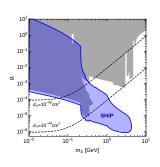
$$\Delta \mathcal{L} = \frac{1}{2} \partial_{\mu} S \partial^{\mu} S + \left(\alpha_{1} S + \alpha S^{2}\right) \left(H^{\dagger} H\right) + \lambda_{2} S^{2} + \lambda_{3} S^{3} + \lambda_{4} S^{4}$$
 | lowest-dim gauge and Lorentz singlet from SM fields | scalar self-couplings |

- $\alpha_1 \neq 0$ : S mixes with Higgs after EW symmetry breaking  $\Rightarrow$  coupling between S and all SM particles
- $\alpha_1 = 0$  (forbidden by exact  $\mathcal{Z}_2$  symmetry): S does not mix with H  $\Rightarrow$ new particles must be pair-produced

## Linear scalar portal







- ightharpoonup Existing limits from searches for rare meson decays e.g. B 
  ightarrow KS
- $\Rightarrow$  Production: K decays (SHiP efficiency  $\approx 0.2\%$ ) and B decays
- ightharpoonup Decay:  $S \to \gamma \gamma, ee, \mu \mu, \pi \pi, KK$



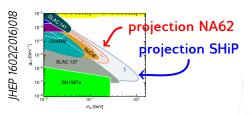
- The SHiP experiment
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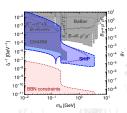
## Axion-like particles





- ightharpoonup The axion mass  $m_A$  is very constrained due to the axial QCD anomaly breaking the PQ symmetry. Other ALPs are not so constrained.
- → SHiP can probe ALPs coupled to gauge bosons and to SM fermions:
  - $pp \to AX, \ A \to \gamma\gamma$ : all neutral, more challenging (left plot)
  - $pp o BX, \ B o AK, \ A o \mu^+\mu^-$  (right plot)







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#### What's next



- crucial to understand the muon flux and spectrum (2018)
  - muon shield design depends on this
  - simulation validated against CHARM measurements, only low- $p_{\perp}$
  - use beam test with a NA61/SHINE-like detector and replica SHiP target
- improve current knowledge of inclusive charm production
  - normalisation for HS signal and  $u_{ au}$  cross-section
  - cascade production? Charmed hadron spectrum?
  - use test beam at H4, with smaller target and nuclear emulsions
  - $10^8$  p.o.t.  $\sim$  4 weeks data taking at the SPS (LS2)



#### Conclusions

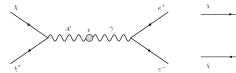
- → General purpose experiment to look for weakly interacting long lived particles
  - covers previously unexplored regions of the Hidden Sector in several theories
  - covers cosmologically interesting regions
- ightharpoonup Unique opportunity for  $u_{ au}$  physics allowing for
  - $ar{
    u}_{ au}$  discovery
  - $\,\sigma$  and form factors measurements
  - also dark matter search
- → Complements LEP/LHC and boosts past experiments sensitivities
  - $\times 10^5$  for HS,  $\times 200$  for  $\nu_{ au}$
  - makes best use of existing SPS complex
- → Next steps:
  - comprehensive design report (2018)
  - $\mu$  flux (2018) and  $\chi_{cc}$  ( $\sim$ 2019) measurements

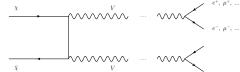




# Motivations for light vector particles





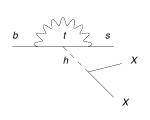


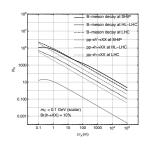
- $\rightarrow$  Dark matter ( $\Omega_{DM} \sim 0.25$ ):
  - light scalar dark matter  $m_\chi \sim {
    m MeV}$  can solve the positron excess
  - WIMP interacting with SM through light mediator ( $\chi \bar{\chi} \to VV \to$  SM) (hides DM from direct searches)
  - non thermal DM (sterile neutrinos)
  - DM self-interaction in structure formation ( $m_V \sim {\sf MeV-GeV}$ )
- $\rightarrow$  Muon g-2:

Light vector particle coupled to muons provides upward correction through one-loop diagram (exchange of A'). *Not* minimal model.

## $\mathcal{Z}_2$ scalar portal





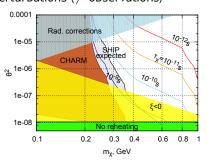


- ightharpoonup Higher dimension portals:  $\frac{1}{\Lambda}|H|^2\psi\psi$  (dark fermions),  $\frac{1}{\Lambda^2}m_{Z_D}^2|H|^2Z_{D\mu}Z_D^{\mu}$  (dark gauge boson)
- → decays of the SM Higgs into hidden states
- → at SHiP  $E_{CM} \simeq 28$  GeV  $< m_H$  Production channels at SHiP:
  - heavy meson decays (dominant is  $B o K^{(*)} XX$ )
  - gluon fusion  $pp \to h^* \to XX$
- → X decays back to SM with different coupling

#### Inflaton



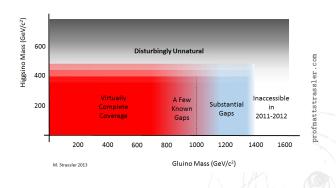
- → In particle physics, the inflaton is a scalar field that couples to SM fields to ensure re-heating of the post-inflation Universe (production of particles that thermalize) and transfer of inflaton fluctuations into adiabatic matter perturbations.
- $ightarrow \, \mathcal{L}_{int} = lpha S^2 H^\dagger H$ , with approx.  $10^{-11} < lpha < 10^{-7}$ 
  - $\alpha < 10^{-11}$   $\longrightarrow$  inefficient reheating
  - $-\alpha > 10^{-7}$   $\longrightarrow$  quantum correction would imply large, scale-dependent density perturbations ( $\neq$  observations)
- Sensitivity at SHiP is dominated by the lifetime exponential.



#### SUSY: where do we stand?



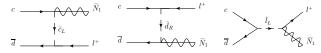
- → SUSY is one of the most popular options to solve naturalness, grand unification and dark matter (WIMP)
- →  $\mathbf{W_{RPC}} = (Y_e)_{ij} L_i H_1 \bar{E}_j + (Y_d)_{ij} Q_i H_1 \bar{D}_j + (Y_u)_{ij} Q_i H_2 \bar{U}_j + \mu H_1 H_2$ SUSY particles produced in pairs. Accelerator searches significantly constrain "natural" scenarios (e.g. MSSM, fine tuning at  $\sim 1\%$ ).



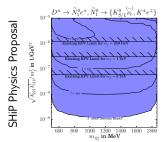
#### SUSY at SHiP: RPV neutralino

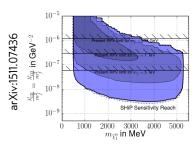


- $\rightarrow \mathbf{W}_{\mathbf{RPV}} = \lambda_{ijk} L_i L_j \bar{E}_k + \lambda'_{ijk} L_i Q_j \bar{D}_k + \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k + \kappa_i L_i H_2$ 
  - The lightest SUSY particle is not anymore stable (no DM)
  - Can be searched for at SHiP in  ${\cal D}$  meson decays:



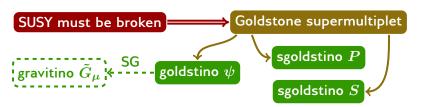
– SHiP sensitivity studied with channels  $\tilde{N}_1^0 o K^0 \overset{(-)}{
u}$  and  $\tilde{N}_1^0 o K^\pm \ell^\mp$ 





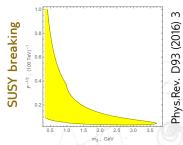
# SUSY at SHiP: sgoldstino





- $\tilde{G}_{\mu}$  ( $\psi$ ) is R-odd
- P, S are R-even  $\Longrightarrow$  can be singly produced and may decay back to pairs of SM particles
- at SHiP:

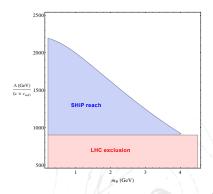
$$\begin{array}{l} pp \xrightarrow{\mathrm{gluon\ fusion}} S \\ D \to SX \\ S \to \ell\ell, \pi\pi \end{array}$$



# SUSY at SHiP: pseudo-Dirac gauginos

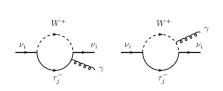


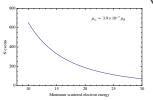
- ightarrow Dirac fermion ( $\Psi$ ) split in two Majorana components ( $\chi_1$ ,  $\chi_2$ )
- interesting dark matter candidate: allows annihilation but appears as Majorana particle for direct and indirect detection purposes
- ightharpoonup Production at SHiP:  $pp o \Psi ar{\Psi}$
- $\rightarrow$  Decay:  $\chi_2 \rightarrow \ell^+ \ell^- \chi_1$



#### Neutrino magnetic moment







If neutrinos are Dirac particles they can get a magnetic moment:

$$\mu_{\nu} = \frac{3eG_F m_{\nu}}{8\pi^2 \sqrt{2}} \simeq (3.2 \times 10^{-19}) \frac{m_{\nu}}{1 \text{ eV}} \mu_B$$

BSM can enhance  $\mu_{\nu}$ .

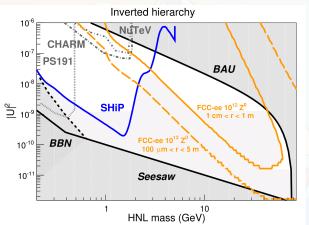
(E.g.: Shrock, Nucl.Phys. B206 (1982) 359)

$$e\nu \to e\nu \Longrightarrow \frac{dN}{dE_e}\Big|_{\mu_{\nu}} = \frac{\pi\alpha^2\mu_{\nu}^2}{m_e^2} \left(\frac{1}{E_e} - \frac{1}{E_{\nu}}\right)$$

Remove BG from  $\nu N$  scattering:  $\theta^2_{\nu e} < 2m_e/E_e \Longrightarrow$  sensitivity:  $N_{evt} \sim 4.3 \times 10^{15} \mu_{\nu}^2/\mu_B^2$ . Prev. limits from  $10^{-7}~(\nu_{\tau})$  to  $10^{-11}~(\nu_e)$ .



## **HNLs at future colliders**



http://arxiv.org/abs/1411.5230 http://arxiv.org/abs/1503.08624

#### Sensitivity with non-zero background



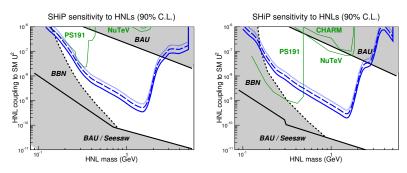
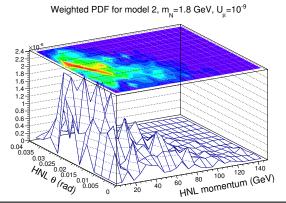


Figure: Variation of the sensitivity contours for scenarios II (left) and IV (right) as a function of the background estimates. The solid blue curve represents the 90% C.L. upper limit assuming 0.1 background events in  $2\times 10^{20}$  proton-target collisions. The dashed blue curve assumes 10 background events. The dotted blue curve assumes a systematic uncertainty of 60% on the level of background, i.e.  $10\pm 6$  background events.

#### Estimating SHiP's physics reach

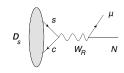
$$\bigotimes$$

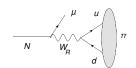
- $\Phi(p.o.t) \times \mathcal{BR}(pp \to NX) \times \mathcal{P}_{vtx} \times \mathcal{BR}(N \to visible) \times \mathcal{A}$ 
  - HNL's momentum and angle are stored in a binned PDF
  - HNL spectra are re-weighted by the probability  $\mathcal{P}_{vtx}(p,\theta\,|\,m_N,U_f^2) \sim \int_V e^{-l/\gamma c \tau} dl$
  - Integral of the weighted PDF gives the total probability  $\mathcal{P}_{vtx}(m_N,U_f^2)$  that HNLs leave a vertex in SHiP's fiducial volume

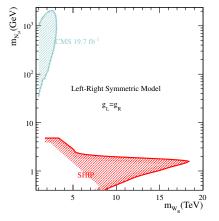


#### Sensitivity in the Left-Right symmetric model









- SHiP limits on  $m_{W_R}$  can be extracted from the HNL limits by  $\left|U_{\mu I}\right|^2 o \left(m_{W_L}/m_{W_R}\right)^4$
- ullet LHC can perform direct searches on both  $W_R$  and  $N_R$
- SHiP can only look for  $N_R$ , but in a domain inaccessible to LHC
- based on CMS, Eur. Phys. J. C 74 (2014) 3149, and Helo, Hirsch, Kovalenko, Phys.Rev. D89 (2014) 073005

#### LFV processes



- ightarrow 
  u oscillations provide evidence of LFV in the neutral sector
- → LFV in charged sector foreseen with  $\mathcal{BR} \sim \mathcal{O}(10^{-40})!$
- $\rightarrow$  New physics models can enhance these  $\mathcal{BR}$ s
  - in seesaw models charged LFV can happen in tree or loop diagrams
  - $\ell \to 3\ell'$  generally favoured with respect to  $\ell \to \ell' \gamma$  (type 2 and 3 seesaw)
- o  $\ell o 3\ell'$  related by unitarity to  $Z^0, h, V o \ell^+\ell'^-$  and  $\ell o \ell'$  conversion in nuclei (most stringent limits so far by SINDRUM II)
  - $au o 3\mu$  and  $\mu o 3e$  can provide better limits than direct searches e.g. for  $\phi o e\mu$ ,  $J/\Psi o e\mu$
  - $\mathcal{BR}( au o 3\mu) < 1.2 imes 10^{-8}$  (BaBar,Belle,LHCb) HFAG, arXiv:1412.7515
- ightharpoonup SHiP will collect  $3\times 10^{15}~ au$  in the forward region
  - requires changes to conceptual design (upgrade):
  - 1 mm W target:  $100 \times$  less  $\tau$ , but decaying outside target
  - LHCb VELO + Si tracker + hadron absorber +  $\mu$  spectrometer
  - sensitivity  $\sim 10^{-10}/\sqrt{N_{\rm targets}}$



## The Hidden Sector

$$L_{world} = L_{SM} + L_{mediation} + L_{HS}$$

- Neutrino portal: new Heavy Neutral Leptons coupling with Yukawa coupling,  $L_{NP}=F_{\alpha I}(\bar{L}_{\alpha}\widetilde{\Phi})N_{I}$
- **Vector portal**: massive dark photon coupling through loops of particles charged both under U(1) and U'(1):  $L_{VP} = \epsilon F'_{\mu\nu} F^{\mu\nu}$
- Scalar portal: light scalar mixing with the Higgs  $L_{SP}=\left(\lambda_iS_i^2+g_iS_i\right)\overline{\Phi}\Phi$
- Axion portal: axion-like particles,  $L_{AP}=rac{A}{4f_A}\epsilon^{\mu\nu\lambda\rho}F_{\mu\nu}F_{\lambda\rho}$
- SUSY: neutralino, sgoldstino, gaugino...

Models	Final states
Neutrino portal, SUSY neutralino	$\ell^{\pm}\pi^{\mp}, \ell^{\pm}K^{\mp}, \ell^{\pm}\rho^{\mp},  \rho^{\pm} \to \pi^{\pm}\pi^{0}$
Vector, scalar, axion portals, SUSY sgoldstino	$\ell^+\ell^-$
Vector, scalar, axion portals, SUSY sgoldstino	$\pi^{+}\pi^{-}, K^{+}K^{-}$
Neutrino portal ,SUSY neutralino, axino	$\ell^+\ell^-\nu$
Axion portal, SUSY sgoldstino	$\gamma\gamma$
SUSY sgoldstino	$\pi^{0}\pi^{0}$

# New Phy

#### New Physics prospects in Hidden Sector



- D = 2: Vector portal
  - Kinetic mixing with massive dark/secluded/paraphoton V :  $\frac{1}{2} \varepsilon F_{\mu\nu}^{SM} F_{HS}^{\mu\nu}$
  - → Motivated in part by idea of "mirror world" restoring left and right symmetry, constituting dark matter, g-2 anomaly, ...
  - Production: proton bremsstrahlung, direct QCD production  $q\bar{q} \to V$ ,  $qg \to Vq$ , meson decays  $(\pi^0, \eta, \ \omega, \ \eta', ...)$
- D = 2: Scalar portal
  - Mass mixing with dark singlet scalar  $\chi:(gS+\lambda S^2)H^\dagger H$
  - → Mass to Higgs boson and right-handed neutrino, inflaton, dark phase transitions BAU, dark matter, "dark naturalness",.
  - Production: Direct  $p + target \rightarrow X + S$ , meson decays e.g.  $B \rightarrow KS$ ,  $K \rightarrow \pi S$
- D = 5/2: Neutrino portal
  - Mixing with right-handed neutrino N (Heavy Neutral Lepton):  $Y_{l\ell}H^{\dagger}\overline{N}_{l}L_{\ell}$
  - → Neutrino oscillation, baryon asymmetry, dark matter
  - Production: Leptonic, semi-leptonic decays of heavy hadrons
- D = 4: Axion portal
  - Mixing with Axion Like Particles, pseudo-scalars pNGB, axial vectors  $a: \frac{a}{\epsilon}G_{\mu\nu}\tilde{G}^{\mu\nu}, \frac{\partial_{\mu}a}{\epsilon}\bar{\psi}\gamma_{\mu}\gamma_{5}\psi$ , etc
  - → Generically light pseudo-scalars arise in spontaneous breaking of approximate symmetries at a high mass scale F
- → Extended Higgs, SUSY breaking, dark matter, possibility of inflaton,...

Production: Primakoff production, mixing with pions and heavy meson decays

- And higher dimensional operator portals
- Chern-Simons portal (vector portal)





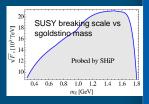




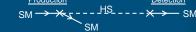
## New Physics prospects in Hidden Sector



- SUper-SYmmetric "portals"
  - Some of SUSY low-energy parameter space open to complementary searches Sgoldstino S(P) :  $\frac{M_{YY}}{F_{UY}}SF^{\mu\nu}F_{\mu\nu}$
  - Neutralino in R-Parity Violating SUSY
  - Hidden Photinos, axinos and saxions....
  - Hidden Photinos, axinos and saxions...



- → A very large variety of models based on these or mixtures thereof
- Two search methods:
  - 1. "Indirect detection" through portals in (missing mass)
  - 2. "Direct detection" through both portals in and out



→ SHiP has significant sensitivity to all of these!

Assumption invisible decay width  $\chi \bar{\chi}$  is absent or sub-dominant,  $m_{\chi} > \frac{1}{2} m_{portal}$ , where  $\chi$  hidden sector particle



## **Sterile Neutrinos**

Fermions get mass via the Yukawa couplings:

$$-\mathcal{L}_{\text{Yukawa}} = Y_{ij}^d \overline{Q_{Li}} \phi D_{Rj} + Y_{ij}^u \overline{Q_{Li}} \tilde{\phi} U_{Rj} + Y_{ij}^\ell \overline{L_{Li}} \phi E_{Rj} + \text{h.c.}.$$

If we want the same coupling for neutrinos, we need right-handed (sterile) neutrinos... the most generic Lagrangian is

$${\cal L}_N = i \overline{N}_i \partial_\mu \gamma^\mu N_i - rac{1}{2} M_{ij} \overline{N^c}_i N_j - Y^
u_{ij} \overline{L_{Li}} ilde{\phi} N_j$$

Kinetic term Majorana mass term Yukawa coupling

Seesaw mechanism:

$$\mathcal{V} = (\nu_{Li}, N_j)$$
  $-\mathcal{L}_{M_{\mathcal{V}}} = \frac{1}{2} \overline{\mathcal{V}} M_{\mathcal{V}} \mathcal{V} + h.c.$   $M_{\nu} = \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix}$   $\lambda_{\pm} = \frac{M_N \pm \sqrt{M_N^2 + 4M_D^2}}{2}$ 

if 
$$M_N \gg M_D$$
:

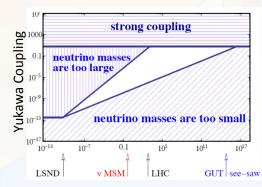
$$\lambda_- \sim \frac{M_D^2}{M_N}$$

$$\lambda_+ \sim M_N$$



## **Sterile neutrino masses**

Seesaw formula 
$$m_D \sim Y_{I\alpha} < \phi >$$
 and  $m_{\nu} = \frac{m_D^2}{M}$ 



- Assuming  $m_{\nu} = 0.1 \text{eV}$
- if  $Y \sim 1$  implies  $M \sim 10^{14} \text{GeV}$
- if  $M_N \sim 1 \text{GeV}$  implies  $Y_{\nu} \sim 10^{-7}$

remember  $Y_{top} \sim 1$ . and  $Y_e \sim 10^{-6}$ 

If we want to explain the smallness of neutrino masses (in a natural way) the mass of sterile neutrinos should be at least at the GeV scale

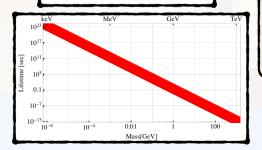
Majorana Mass (GeV)



# Constraints on N<sub>1</sub>

The decay mode N o 
u 
u 
u is always present

$$LT = \left(\frac{U^2 G_F^2 M_N^5}{86 \pi^3}\right)^{-1} \simeq 0.3 \left(\frac{1 GeV}{M_N}\right)^4 sec$$



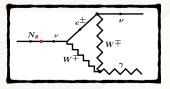
This gives an upper bound for the mass of the mass of the sterile neutrino Dark Matter

- $M_N \sim 1 KeV \Longrightarrow \tau_N \sim 10^{24} sec$
- Age of the Universe  $\sim 10^{-6}$

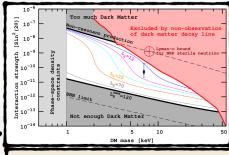


# Constraints on N<sub>1</sub>

DM sterile neutrinos decay subdominantly as  $N_1 \to \nu \gamma$  with a branching ration  $\mathcal{B}(N_1 \to \gamma \nu) \sim \frac{1}{123}$ 



Discussion in the community, not yet clear if this is a "good" signal, needs confirmation



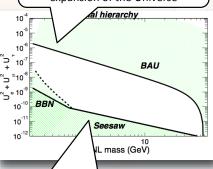
Bulbul et al. 2014 (arXiv:1402.2301)

Boyarsky et al. 2014 (arXiv:1402.4119)



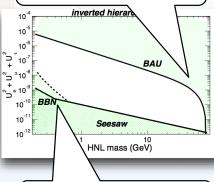
# Constraints on N<sub>2</sub>, N<sub>3</sub>

If  $U^2$  is too large,  $N_{2,3}$  are in **thermal equilibrium** during the expansion of the Universe



The **seesaw** limit defines the region where  $N_{2,3}$  can explain the observed active neutrino  $\Delta m^2$ 

At  $M_N \ge M_W$  the rate is **enhanced** by  $N \to Wl$  leading to stronger constraints on  $U^2$ 



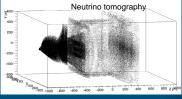
If  $\tau(N_2, N_3) < 0.1 \, s$ , they cannot affect the **Big Bang nucleosynthesis** 



#### Backgrounds with TP detector



Background source	Decay modes
$\nu$ or $\mu$ + nucleon $\rightarrow X + K_L$	$K_L \rightarrow \pi e \nu, \pi \mu \nu, \pi^+ \pi^-, \pi^+ \pi^- \pi^0$
$\nu$ or $\mu$ + nucleon $\rightarrow X$ ] + $K_S$	$K_S \rightarrow \pi^0 \pi^0, \pi^+ \pi^-$
$\nu$ or $\mu$ + nucleon $\rightarrow X + \Lambda$	$\Lambda \rightarrow p\pi^-$
$n \text{ or } n + \text{nucleon} \rightarrow X + K_T \text{ etc.}$	as above

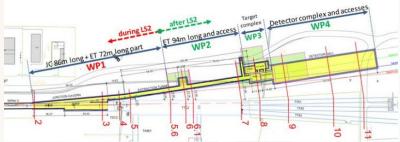


- Background summary: no evidence for any irreducible background
  - No events selected in MC → Expected background UL @ 90% CL

Background source	Stat. weight	Expected background (UL 90% CL)	
$\nu$ -induced		1 0 ( /	
$2.0$	1.4	1.6	
$4.0$	2.5	0.9	
p > 10  GeV/c	3.0	0.8	
$\overline{\nu}$ -induced			
$2.0$	2.4	1.0	
$4.0$	2.8	0.8	
p > 10  GeV/c	6.8	0.3	
Muon inelastic	0.5	4.6	
Muon combinatorial	-	< 0.1	
Cosmics			
p < 100  GeV/c	2.0	1.2	
p > 100  GeV/c	1600	0.002	



# NA work packages



- Preparation of facility in four well-defined quasi-independent work packages
  - WP1: Junction cavern + 70m beam line for clearance during operation (21 months)
  - WP2 : Rest of beam line (12 months)
  - WP3 : Target complex (12 months)
  - · WP4 : Experiment facility (18 months)
  - → Only WP1 has to be done during a stop of the North Area only
  - → WP1 associated with cool down, removal and re-installation of services and beam line (24-27 months)
  - → Construction of facility has no interference with operation of SPS and LHC at any time



#### SHiP target



#### Design considerations with 4x10<sup>13</sup> p / 7s

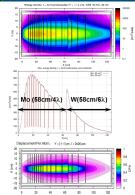
- → 355 kW average, 2.56 MW during 1s spill
  - · High temperature
- Compressive stresses
- Atomic displacement
- Erosion/corrosion
- Material properties as a function of irradiation
- Remote handling (Initial dose rate of 50 Sv/h...)
- → Hybrid solution: Mo allow TZM  $(4\lambda)$  + W $(6\lambda)$

	DONUT 1)	CHARM 2)	SHiP
Target material	W-alloy	Cu (variable $\rho$ )	TZM + pure W
Momentum (GeV/c)	800	400	400
Intensity	0.8*1013	1.3*10 <sup>13</sup>	4*10 <sup>13</sup>
Pulse length (s)	20	23*10-6	1
Rep. rate (s)	60	~10	7.2
Beam energy (kJ)	1020	830	2560
Avg. beam power (spill) (kW)	51	3.4*10 <sup>7</sup> (fast)	2560
Avg. beam power (SC) (kW)	17	69	355
POT	Few 10 <sup>17</sup>	Few 10 <sup>18</sup>	2*10 <sup>20</sup>



30x30cm<sup>2</sup>

**TZM** 





-2

#### Active muon shield

CÉRN

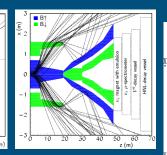
2800 tonnes

48m

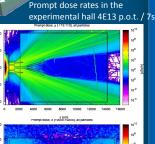
- Muon flux limit driven by emulsion based  $\nu$ -detector and "hidden particle" background
- Passive and magnet sweeper/passive absorber options studied:
  - Conclusion: Shield based entirely on magnetic sweeping with  $\int B_{\gamma} \, dl \sim$  86 Tm
  - $\rightarrow$  <7x10<sup>3</sup> muons / spill (E<sub>u</sub> > 3 GeV) which can potentially produce V0 (K<sub>L</sub>)
  - → Negligible occupancy

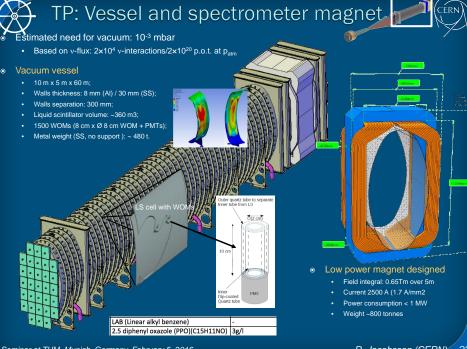
350 GeV

+10 mrad 0 mrad -10 mrad



→ Challenges: flux leakage, constant field profile, modelling magnet shape





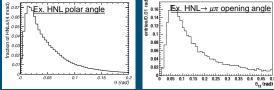


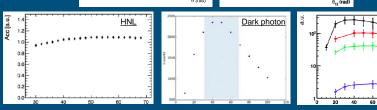
#### HS detector optimization





- Optimization of geometrical acceptance for a given  $\mathsf{E}_{\mathsf{beam}}$  and  $\Phi_{\mathsf{beam}}$ 
  - Hidden particle lifetime (~flat for longlived)
  - Hidden particle production angles (~distance and transversal size)
  - Hidden particle decay opening angle (~length and transversal size)
  - Muon flux (~distance and acceptable occupancy)
  - Background (~detector time and spatial resolution)
  - Evacuation in decay volume / technically feasible size ~ W:5m x H:10m





→ Acceptance saturates ~40m - 50m

Hidden scalar

#### HS tracking system

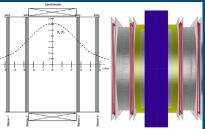
#### A62-like straw detector

Parameter	Value
Straw	
Length of a straw	5 m
Outer straw diameter	9.83 mm
Straw wall (PET, Cu, Au)	
PET foil thickness	$36 \mu m$
Cu coating thickness	50 nm
Au coating thickness	20 nm
Wire (Au-plated Tungsten)	
diameter	$30 \mu m$
Straw arrangement	
Number of straws in one layer	568
Number of layers per plane	2
Straw pitch in one layer	17.6 mm
Y extent of one plane	$\sim 10 \text{ m}$
Y offset between straws of layer 1&2	8.8 mm
Z shift from layer 1 to 2	11 mm
Number of planes per view	2
Y offset between plane 1&2	4.4 mm
Z shift from plane 1 to 2	26 mm
Z shift from view to view	100 mm
Straw station	
Number of views per station	4 (Y-U-V-Y)
Stereo angle of layers in a view Y,U,V	0, 5, -5 degrees
Z envelope of one station	$\sim 34~\mathrm{cm}$
Number of straws in one station	9088
Straw tracker	
Number of stations	4
Z shift from station 1 to 2 (3 to 4)	2 m
Z shift from station 2 to 3	5 m
Number of straws in total	36352

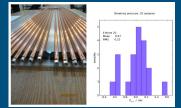
#### Straws in test beam 2016

- · Study sagging effects and compensation
- Read out of signal, attenuation / two-sided readout
- Upstream straw veto may be based on same technology

#### Horizontal orientation of 5m straws



#### First production of 5m straws at JINR

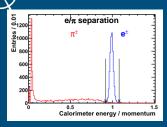


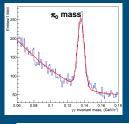
JINR Dubna (NA62, SHiP): Straws St Petersburg (CMS, SHiP): Infra

## PID performance



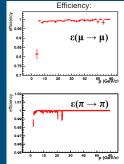


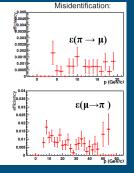




Electron efficiency >98%
Pion contamination:<2%
Neutral pion mass resolution: 5 MeV

Muon misid with ECAL+HCAL Rejection factor for  $\varepsilon_u$ =95% Energy, GeV E+H1+H2 1.0 23 1.5 32 2.0 50 2.7 120 3.0 160 5.0 210 2/07/1000 250





→ ECAL (July), HCAL (September), MUON (October) in test beam 2015 on PS and SPS